## BRIEF COMMUNICATIONS

# The Series of Spinels $\mathrm{Co}_{3-s} \mathrm{Al}_{s} \mathrm{O}_{\mathbf{4}}(\mathbf{0}<\boldsymbol{s}<\mathbf{2})$ : Study of $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ 

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#### Abstract

Co (III) of $\mathrm{CO}_{3} \mathrm{O}_{4}$ can be gradually replaced by Al to produce the series of spinels $\mathrm{Co}_{3-5} \mathrm{Al}_{5} \mathrm{O}_{4}(0<s<$ 2), S.G. Fd $3 m$ (No. 227), $Z=8$. For $\mathrm{Co}_{2} \mathrm{AlO}_{4}\left(s=1\right.$ ), $a=8.086(1) \AA, U=528.7(2) \AA^{3}, D_{x}=5.25 \mathrm{Mgm}^{-3}$, $u=0.264$, and $27 \%$ of Al in $8(a)$ positions, $R=0.031$. The frequencies of the observed ir absorption bands of $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ are also presented.


## Introduction

The spinels are crystalline solids well known to mineralogists, chemists, and solid-state physicists. The most abundant spinel groups, of stoichiometry $A B_{2} X_{4}$, are those formed by some metals in oxidation states II-III and II-IV. $X$ usually stands for oxygen or one chalcogen. Stoichiometries as $\mathrm{Co}_{7 / 3} \mathrm{Sb}_{23} \mathrm{O}_{4}, \mathrm{Li}_{4 / 3} \mathrm{Ti}_{5 / 3} \mathrm{O}_{4}, \mathrm{Cu}_{1 / 2} \mathrm{Fe}_{52} \mathrm{O}_{4}$, $\mathrm{LiAlSiO}_{4}, \mathrm{LiZnSbO}_{4}, \mathrm{CrMnNiO}_{4}$, are also known. All these compounds can be represented by $M_{3} X_{4}, M$ denoting the same or different metals. For a nondeficient spinel containing two metals, $M, M^{\prime}$, in oxidation states $m, m^{\prime}$, respectively, and divalent $X$ anions, the formula $M_{\mathrm{b}} M_{\mathrm{c}}^{\prime} X_{4}$, wherein $b=$ $\left(3 m^{\prime}-8\right) /\left(m^{\prime}-m\right) ; c=(8-3 m) /\left(m^{\prime}-\right.$ $m$ ); includes all the possible stoichiometries, i.e., $\quad M_{2}^{1} M^{\mathrm{VI}^{\mathrm{I}} X_{4}} ; \quad M{ }_{1 / 4} M_{5 / 4}^{\mathrm{Y}} X_{4}$;

[^0]$M_{4 / 3}^{\mathrm{L}_{3}} M_{53}^{\mathrm{V}} X_{4} ; \quad M_{1_{2}}^{\mathrm{I}_{2}} M_{51}^{\mathrm{II}} X_{4} ; \quad M_{5_{2}}^{\mathrm{I}} M_{1 / 2}^{\mathrm{Y}} X_{4} ;$ $M_{73}^{\mathrm{H}} M_{23}^{\mathrm{V}_{3}} X_{4} ; M_{2}^{\mathrm{II}} M^{\mathrm{IV}} X_{4} ; M^{\mathrm{I}} M_{2}^{\mathrm{II}} X_{4}$.
$b$ and $c$ values are obtained if electroneutrality is assumed and by requiring that the sum of $b$ and $c$ has to be equal to three. For three metals, $M, M^{\prime}, M^{\prime \prime}$, in oxidation states $m, m^{\prime}, m^{\prime \prime}$, the series $M_{\mathrm{s}} M_{\mathrm{d}}^{\prime} M_{\mathrm{e}}^{\prime \prime} X_{4}$, where $d=\left(8-3 m^{\prime \prime}-m s+m^{\prime \prime} s\right) /\left(m^{\prime}-\right.$ $\left.m^{\prime \prime}\right) ; \boldsymbol{e}=\left(3 m^{\prime}-8+m s-m^{\prime} s\right) /\left(m^{\prime}-m^{\prime \prime}\right)$ is encountered with $s, d$, and $e$ positive.
The spinel $\mathrm{CoAl}_{2} \mathrm{O}_{4}$, the so-called Thénard blue, has been known for a long time and has been extensively studied ( 1 , 2). It may be considered as derived from $\mathrm{Co}^{\mathrm{II}} \mathrm{CO}_{2} \mathrm{IIH}_{4}$ by replacing all the $\mathrm{Co}^{\mathrm{III}}$ with Al. This paper deals with the $\mathrm{Co}^{\mathrm{II}}$ $\mathrm{Co}_{2}^{\mathrm{HI}} \mathrm{Al}_{s} \mathrm{O}_{4}(0<s<2)$ series, which arises when the above-mentioned substitution is carried out in steps. The spinel structure type, $M_{3} X_{4}$, is based on a cubic closestpacked array of $X^{2-}$ ions, with the $M$ ions occupying one-eighth of the tetrahedral and
one-half of the octahedral sites, $M_{x}\left[M_{3-x}\right]$. The brackets represent octahedral site occupations and $M$ may represent as before, either one or several metals. In a spincl $M M_{2}^{\prime} X_{4}$ two extreme distributions of the cations among the available sites are possible: the so-called normal, $M\left[M_{2}^{\prime}\right] X_{4}$, and the inverse distribution, $M^{\prime}\left[M M^{\prime}\right] X_{4}$. Intermediate cation distributions, $M_{1-x} M_{x}^{\prime}$ [ $M_{x} M_{2-x}^{\prime}$ ], are also possible. Here $x$ is the degree of inversion, equal to zero and unity for the normal and inverse arrangements, respectively. When referred to space group $F d 3 m, Z=8$, and origin at center, $(\overline{3} m)$, the 24 cations occupy 8 tetrahedral (a) positions, $1 / 8,1 / 8,1 / 8$, and 16 octahedral (d) positions, $1 / 2,1 / 2,1 / 2$; and the anions are located on $32(e)$ positions, $u, u, u$. Therefore, letting $a$ be the unit-cell parameter, the significant distances in the spinel unitcell are
At tetrahedral (a) positions:
\[

$$
\begin{aligned}
& M-X=a 3^{1 / 2}(u-1 / 8) \\
& X-X=a 2^{1 / 2}(2 u-1 / 4)
\end{aligned}
$$
\]

At octahedral (d) positions:

$$
\begin{aligned}
& M-X=a\left(3 u^{2}-2 u+3 / 8\right)^{1 / 2} \\
& X-X=a 2^{1 / 2}(3 / 4-2 u) \\
& X-X^{\prime}=a\left(4 u^{2}-2 u+3 / 8\right)^{1 / 2} .
\end{aligned}
$$

Variations in $u$ reflect the adjustment of the structure to accommodate differences in the relative radii of the cations in the tetrahedral and octahedral sites. When $u=$ 0.250 , the $X$ anions are arranged in an ideal cubic closest-packing arrangement. This particular situation and the cases $u<0.250$ are rarely realized; more often $u$ value lies in the range 0.251 to 0.269 .

In order to study a particular spinel, three variables have to be established: the cell dimension, $a$; the positional parameter, $u$, for oxygen atoms; and the distribution of cations, $x$, between tetrahedral and octahedral sites. It is the aim of this paper to establish these parameters for $\mathrm{Co}^{\mathrm{II}} \mathrm{Co}^{\mathrm{III}} \mathrm{AlO}_{4}$, middle term of the above-mentioned spinel series.

## Experimental

Mixtures of analytical grade $\mathrm{Co}_{3} \mathrm{O}_{4}$ and $\gamma$ $\mathrm{Al}_{2} \mathrm{O}_{3}$ in the molar ratios $\mathrm{Co}_{3} \mathrm{O}_{4}: \mathrm{Al}_{2} \mathrm{O}_{3}=(6$ $-2 s): 3 s(s=0.25,0.50 \ldots 1.75)$, were heated in air at 1073 K for 24 hr , to obtain seven samples of the series $\mathrm{Co}_{3-s} \mathrm{Al}_{s} \mathrm{O}_{4}$. Experimental conditions and details on the analysis of the X-ray diffraction data have been described elsewhere (3). The unit-cell parameter of $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ was refined from the $2 \theta$ values of the last eight reflections; the arbitrarily chosen Debye-Waller factors, $B_{j}$, for the calculation of the intensities, were 0.80 and $0.60 \AA^{2}$ for the oxygen and Co atoms, respectively; the value, $0.75 \AA^{2}$, for Al was obtained from these two by interpolation of the atomic mass of Al between those of Co and O . The intensities were calculated in two stages. First, six runs including 36 jobs were performed for $u$ $=0.250,0.255,0.260,0.2625,0.265,0.270$ and $x=0,10,20,30,33,40$. The second stage included 54 jobs for the following $u$ and $x$ values: $u=0.261,0.262,0.2625$, $0.263,0.264,0.265,0.266,0.267,0.268 ; x=$ $25,26,27,28,29,30$. Infrared spectra were recorded using a Perkin-Elmer 325 spectrophotometer, KBr disks and Nujol oil mulls.

## Results

The seven samples were obtained as blackish blue powders, which produced good X-ray diffraction patterns characteristic of spinels, space group Fd3m (No. 227). The $a$ values exhibited slight variations, confirming the existence of the $\mathrm{Co}^{\mathrm{II}}$ $\mathrm{Co}_{2-s}^{\mathrm{IH}} \mathrm{Al}_{s} \mathrm{O}_{4}(0<s<2)$ series. The observed and calculated interplanar spacings, $d_{0}$ and $d_{\mathrm{c}}$, for $\mathrm{CO}_{2} \mathrm{AlO}_{4}$, the middle term of the series, are included in Table I. Thus, $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ is a spinel, with $a=8.086(1) \AA, U$ $=528.7(2) \AA^{3}, Z=8$, and $D_{x}=5.25 \mathrm{Mgm}^{-3}$. The observed, $I_{\mathrm{o}}$, and calculated, $I_{\mathrm{c}}$, inten-

TABLE I
X-Ray Powder Diffraction Data for $\mathrm{Co}_{2} \mathrm{AlO}_{4}$

| $h k l$ | $d_{0}$ | $d_{c}$ | $I_{0}$ | $I_{c}$ |
| :---: | :---: | :---: | ---: | ---: |
| 111 | 4.670 | 4.669 | 147 | 149 |
| 220 | 2.860 | 2.859 | 352 | 348 |
| 311 | 2.439 | 2.438 | 1,052 | 1,000 |
| 222 | 2.3337 | 2.3342 | 64 | 48 |
| 400 | 2.0214 | 2.0215 | 207 | 235 |
| 422 | 1.6505 | 1.6505 | 95 | 96 |
| 511,333 | 1.5568 | 1.5562 | 356 | 362 |
| 440 | 1.4294 | 1.4294 | 420 | 431 |
| 531 | 1.3668 | 1.3668 | 10 | 13 |
| 600,442 | - | 1.3477 | - | 1 |
| 620 | 1.2784 | 1.2785 | 25 | 30 |
| 533 | 1.2331 | 1.2331 | 86 | 84 |
| 622 | 1.2189 | 1.2190 | 33 | 33 |
| 444 | 1.1672 | 1.1671 | 23 | 26 |
| 711,551 | - | 1.1323 | - | 10 |
| 642 | 1.0806 | 1.0805 | 34 | 45 |
| 731,553 | 1.0528 | 1.0527 | 124 | 124 |
| 800 | 1.0107 | 1.0108 | 48 | 52 |
| 733 | - | 0.9879 | - | 1 |
| 822,660 | 0.9530 | 0.9529 | 24 | 25 |
| 751,555 | 0.9336 | 0.9337 | 92 | 91 |
| 662 | 0.9275 | 0.9275 | 24 | 20 |
| 840 | 0.9041 | 0.9040 | 35 | 34 |
| 911,753 | - | 0.8876 | - | 6 |
| 842 | - | 0.8823 | - | 1 |
| 664 | - | 0.8620 | - | 9 |
| 931 | 0.8476 | 0.8476 | 104 | 106 |
| 844 | 0.8254 | 0.8253 | 219 | 201 |
| $933,771,755$ | - | 0.8127 | - | 5 |
| $10.20,862$ | 0.7930 | 0.7929 | 66 | 70 |
| 951,773 | 0.7817 | 0.7817 | 306 | 326 |
|  |  |  |  |  |

sities which gave the best discrepancy factor, $R=\left(\Sigma\left|I_{o}^{1 / 2}-I_{\mathrm{c}}^{1 / 2}\right|\right) / \Sigma I_{\mathrm{o}}^{1 / 2}=0.031$, are also included in Table I, and were obtained for $u=0.264$, with the origin at center, ( $\overline{3} m$ ), and $x=27 \%$ of Al atoms in tetrahedral ( $a$ ) sites. The M-O and O-O distances for $\mathrm{Al}_{0.27} \mathrm{Co}_{0.73}\left[\mathrm{Al}_{0.73} \mathrm{Co}_{1.27}\right] \mathrm{O}_{4}$ were: for tetrahedral (a) sites, $\mathrm{M}-\mathrm{O}=1.947 \AA, \mathrm{O}-\mathrm{O}=$ $3.179 \AA$; and for octahedral (d) sites, M-O $=1.915 \AA, \mathrm{O}-\mathrm{O}=2.539 \AA, \mathrm{O}-\mathrm{O}^{\prime}=2.868$ $\AA$.

Table II includes the frequencies $\left(\mathrm{cm}^{-1}\right)$ of the observed infrared absorption bands

TABLE II
Frequencies ( $\mathrm{cm}^{-1}$ ) of the ir Absorption Bands Observed for $\mathrm{Co}_{2} \mathrm{AlO}_{4}$, as Compared with those

Reported ${ }^{a}$ For $\mathrm{Co}_{3} \mathrm{O}_{4}$ and $\mathrm{CoAl}_{2} \mathrm{O}_{4}$

| Compound | $\nu_{1}$ | $\nu_{2}$ | $\nu_{3}$ | $\nu_{4}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ | 656 | 564 | 383 | 217 |
| $\mathrm{Co}_{3} \mathrm{O}_{4}$ | 672 | 590 | 392 | 220 |
| $\mathrm{CoAl}_{2} \mathrm{O}_{4}$ | 684 | 563 | 523 | 237 |
| ${ }^{a} \mathrm{~J}$. Preudhomme and P. Tarte, Ref. (4). |  |  |  |  |

for $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ as well as those for $\mathrm{Co}_{3} \mathrm{O}_{4}$ and $\mathrm{CoAl}_{2} \mathrm{O}_{4}$ measured by Preudhomme and Tarte (4). These authors studied the infrared spectra of normal $A^{\mathrm{II}}\left[B_{2}^{\mathrm{II}}\right]$ spinels $(x=$ 0 ) in which four bands, corresponding to the four infrared-active modes $T_{1 u}$, are observed. They showed that the two $\nu_{1}$ and $\nu_{2}$ bands are related to the vibrations of the lattice of octahedral groups, and that two low-frequency bands, $\nu_{3}$ and $\nu_{4}$, correspond to complex vibrations involving the participation of both the tetrahedral and the octahedral cations. These assignments were deduced for the normal cation distributions and cannot be applied to inverse spinels. More recently, Toriumi and co-workers (5) have shown that $\mathrm{CoAl}_{2} \mathrm{O}_{4}$ is not a normal spinel ( $x=0.155$ ). The assignments proposed by Preudhomme and Tarte (4) thus are not applicable to $\mathrm{CoAl}_{2} \mathrm{O}_{4}$. This is also the case for $\mathrm{Co}_{2} \mathrm{AlO}_{4}$. Nevertheless, the partial substitution of $\mathrm{Co}^{\mathrm{III}}$ from $\mathrm{CO}_{3} \mathrm{O}_{4}$ by Al to give $\mathrm{Co}_{2} \mathrm{AlO}_{4}$ is accompanied by a rather slight lowering of the four frequencies. By contrast, $\mathrm{CoAl}_{2} \mathrm{O}_{4}$ exhibits an erratic variation of these frequencies which is especially pronounced for $\nu_{3}$.

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